

QUICK GUIDE to QSM(Quantum Stochastic Mechanics with only Axiom of Quantization Principle) :

① : Serious Difficulties of Quantum Mechanics lie in "Time and Hamiltonian Structure" :

None can succeed to establish non-equilibrium statistical mechanics as quantum principle, which is essentially due to lack of genuine recognition on "time". Time in quantum mechanics(=QM) is non-observable, but is quantum statistical variable like as temperature. In QM, physical variable's very important "hermiteness" is equivalent to physical observability. The difficulty of time in QM is easily proved being nothing quantum state transition under "observable hermitian Hamiltonian $\equiv H_0$ ". That is, $i\hbar\partial_t\Psi = H_0\Psi$ (=S eqn) is entirely stationary. By the way, QM structure can be constructed from canonical quantization on conjugate variables as $[P_j, Q_j] \equiv P_j Q_j - Q_j P_j = i\hbar 1$ or $\Delta P \Delta Q \geq 1/2\hbar$, in which dimension $[\hbar] = \text{time} \times \text{energy} = [\text{action}]$. Hence conjugate variable of time "t" is observable Hamiltonian H_0 in which "t" can not be observable. In order to realize quantum state transition, we can not help to introduce non-hermitian Hamiltonian $\equiv H_s(t)$ where $\Delta E = \infty$ due to "non-observability of energy" and also $\Delta t = \hbar/\Delta E = 0$. Phenomena of $H_s(t)$ must be instantaneous quantum state transition and agree with physical realities. Schroedinger eqn is generally derived from Quantization Principle(CP) especially as for time and energy. "QP" can construct any feature of quantum mechanics except interaction force problems.

② : Hermite Hamiltonian's Complete Causality and Markovian of General Quantum Process :

As H_0 is entirely mathematically regular(=analytical), so H_0 's phenomena must realize entirely causal uniqueness. This fact can be proved as "unique eigen state realization of H_0 's maximum observables P_j " in which $[P_j, H_0] = 0$ by using H_0 's stationarity $\Delta t = \infty$. Then paradox of Shrodinger's dog is completely resolved and establish "quantum process'es complete Markovian feature". That is, "a general quantum sample process becomes series of transitions among H_0 's unique eigen state". In the matter of course, such statistical transitions are caused by $H_s(t)$'s series of the realization. $H_s(t)$ is non-causal (statistical quantum transition) and is due to "its mathematical singularity" which is entirely reasonable due to Coedel incompleteness theorem and justify being of divergence difficulty in perturbation integral. Hence "renormalization method" becomes final answer in standard theory.

③ : Stochastic Hamiltonian and its Statistics (Evolution Theorem(or Principle) by Energy Fluctuation) :

After all, in realities, time dependent Hamiltonian becomes random alternating realization of H_0 & $H_s(t)$ like as that of "long term conservative regime $\langle H_0 \rangle$ and short time revolutionary one $\langle H_s(t) \rangle$ ", so time dependent Schroedinger equation must be stochastic differential one with stochastic Hamiltonian $H_R(t)$. Hence we must establish statistics of $H_s(t)$, which can be derived as probability density function $\Theta(t) = \Delta E(t)/\hbar$ for realizing spot time duration of $H_s(t)$ on time axis from modified Winer-Kintchin Theorem. $\Delta E(t)$ is statistical deviation of energy in statistical ensemble (Evolution Theorem by Energy Fluctuation $\equiv \text{ETEF}$). Thus necessary tools are almost derived for establishing "Quantum Stochastic Mechanics".

④ : Establishing Quantum Stochastic Mechanics and the Problems of General Isolated Closed Systems :

Once Markovian had been proved, there must be Master Equation as conservation law of probability flow for state transitions. Now $\omega_{jk}(t)$ = state j's density, $\Gamma_{jk}(t)$ = state transition $k \rightarrow j$ probability rate. Then $\partial_t \omega_j(t) = \sum_k \Gamma_{jk}(t) \omega_k(t) - \sum_k \Gamma_{kj}(t) \omega_j(t) \equiv \{\text{inflow to } |j\rangle \text{ from } |k\rangle\} - \{\text{outflow from } |j\rangle \text{ to } |k\rangle\}$. $\Gamma_{jk}(t) = \Theta(t) T_{jk} = [\Delta E(t)/\hbar] T_{jk}$ is derived, where T_{jk} is 1st order transition probability by $H_s(t)$. Finally we derive QSM Master Eqn representing isolated closed thermo-system in statistical ensemble.

$$(1) \partial_t \omega_j(t) = \langle \Delta E(t)/\hbar \rangle \sum_k [T_{jk} - \delta_{jk}] \omega_k(t). \quad \langle \text{reaction rate} \equiv 1/\Delta t(t) = \Delta E(t)/\hbar : \text{uncertainty of } \Delta E \text{ \& } \Delta t \rangle$$

The equation yields following important realities on isolated closed thermo-dynamical system. (2) irreversibility of the Eqn, (3) entropy increasing law, (4) general relaxation process solution which means that any "thermo-chemical reactions" in isolated closed system shall stop their reactions at last.

⑤ : The Problems of General Opened Systems with Thermo-Chemical External Flows :

The method was generalized to opened system with thermo-chemical external flows in the equation (5). Then we derived also (6) "heat beating solution in constant flows" by mathematically simplified modeling.

$$(5) \partial_t \omega_j(t) = \langle \Delta E_R(t)/\hbar \rangle \sum_k [T_{jk} - \delta_{jk}] \omega_k(t) + \langle \Delta E_L(t)/\hbar \rangle \sum_k [L_{jk} - \delta_{jk}] \omega_k(t)$$

where $\langle \Delta E_L(t)/\hbar \rangle L_{jk}$ is state transition $k \rightarrow j$ probability rate caused by "external quantum flow".

Note on Goedel's Completeness Theorem(≡CT)(1928) and Incompleteness Theorem(≡IT)(1930) :

CT: Any deterministic true proposition T in in contradiction theory K (≡ axiom system) is provable.

IT: There must be also indeterministic proposition X in in contradiction theory K with $N^{\infty}T$.

☞: Author proved that generally X is probabilistical phenomena caused from singularity in K. So called "chaos" is mere a deterministic sample process of stochastic ensemble. Therefore any proposition must be either deterministic or statistical (≡ The Ultra Completeness Theorem ≡ UCT). Now X in the natural number theory N (≡ $N^{\infty}T$) is "the maximum number in N". That is infinity $\equiv \infty$. Then real number zero $\equiv 0^* = 1/\infty$. Therefore 0^* is also indeterministic! (∞ and 0^* are origin of mathematical singularity).

○: Quantization Principle on Canonical Conjugate Variable as Fundamental Axiom of Quantum Mechanics :

①: Quantization :

Generally fundamental dynamical structure of Quantum Mechanics (QM) and the standard Quantum Field Theory (QFT) are constituted from Quantization Principle on canonical conjugate variables. $L = L(Q_i, \dot{Q}_i)$ is Lagrangean, then conjugate variable P_j of Q_i is defined as $P_j \equiv \partial L / \partial (\dot{Q}_j)$. Then their quantization is as follows. These algebra define them as quantum numbers or operators.

(1): $[P_j, Q_k] = i\hbar \delta_{jk}$, (2): $[P_j, P_k] = [Q_i, Q_k] = 0$.

☞: Relich-Dixmier theorem: such variables $[P_j, Q_i] = i\hbar$ are transformed into $\{P_j = -i\hbar \partial / \partial x_j, Q_i = x_j\}$ by certain unitary transform in general.

②: Deriving Schroedinger Equation :

As for time "t", its has two canonical conjugate variable as $\{H_0 \equiv \text{Hamiltonian}, i\hbar \partial_t \equiv \text{time derivative}\}$. Because $x_\mu = (ict, x_1, x_2, x_3)$ and $p_\mu = \{iE/c, p_1, p_2, p_3\} = -i\hbar \partial_\mu$. Hence $E = (c/i)(-i\hbar \partial / \partial x_0) = (c/i)(-i\hbar \partial / \partial ict) = i\hbar \partial / \partial t$. Then $H_0 = i\hbar \partial_t$ for any functions? Absolutely no. Unique possibility is $H_0 = i\hbar \partial_t$ for certain function Ψ . That is $i\hbar \partial_t \Psi = H_0 \Psi, \dots (1)$

③: Uncertainty Theorem for Canonical Conjugate Variable $[P_j, Q_i] = i\hbar$:

Then we can prove following inequality in which $\Delta P_j, \Delta Q_i$ are statistical deviation.

$\Delta P_j \Delta Q_i \geq \frac{1}{2} \hbar \dots (1)$

☞: As for time and energy, it is rather complicated to derive. The result is $\Delta E \Delta t = \hbar, \dots (2)$.

(2) is derived from modified Winer-Kintchin theorem as evolution theorem by energy fluctuation.

Then ΔE is statistical deviation of concerned system and Δt is average reaction time. See \rightarrow ⑤.

●: Serious Difficulties of Quantum Mechanics lie in "Time and Hamiltonian Structure" :

① Observable Hermitian Operator's Spectral Representation :

(1) General linear operator A in function space: $\{|r\rangle \equiv \text{orthogonal function set}, \langle r|s\rangle = \delta(r-s)$

$A |r\rangle = \int ds a_{s,r} |s\rangle = \int ds \int dt a_{s,t} |s\rangle \langle t|r\rangle. \rightarrow A = \int ds \int dt a_{s,t} |s\rangle \langle t|$.

(2) General hermite operator :

$\langle s| A |r\rangle = \langle r| A |s\rangle^* = a_{s,r} = a_{r,s}^*$. (diagonalizationability by Unitary transform U).

(3) spectral representation of hermite operator: $\langle (\int ds U_p a_{s,t}) \equiv A_p \delta(p-t) : A_p = \text{real number}$

$A^{\dagger} \equiv U \cdot A = \int dp \int dq U_{p,q} |p\rangle \langle q| \times \int ds \int dt a_{s,t} |s\rangle \langle t| = \int dp \int dt \int ds U_{p,s} a_{s,t} |p\rangle \langle t| = \int dp \int dt A_p \delta(p-t) |p\rangle \langle t| = \int dp A_p |p\rangle \langle p|$.

② $H_0 \equiv$ hermitian Hamiltonian's Stationality and the Causalitical Uniqueness :

After all, simultaneous determination on "time and energy" is impossible due to uncertainty of 0 ③(2). Therefore H(t) of time dependent Hamiltonian never can be deterministic one. Thus causalitical form of time dependent theory becomes impossible in quantum theory in general. Unfortunately this serious facts has been neglected to cause serious difficulties in various aspects of quantum physics.

Nothing quantum state transition by hermitian Hamiltonian H_0 . That is, absolutely stationary! And also under hermite Hamiltonian, realizable quantum state must be unique eigen one in its maximum observable. So called "superpositional state of H_0 's eigen ones are forbidden in general". Hence famous paradox of Schrodinger dog is resolved. Though such as for position variable of non maximum observable of H_0 , superpositional state of the variable's eigen state is realized. After all, the uniqueness feature ensures general quantum processes Markovian nature and enable establishing Quantum Stochastic Mechanics.

(1) "Nothing Quantum State Transition Under Hermitian Hamiltonian $\equiv H_0$:

$$H_0 = \int d\epsilon |\epsilon\rangle\langle\epsilon| ; \psi = \int d\epsilon a(t;\epsilon) |\epsilon\rangle. \rightarrow i\hbar \partial_t \psi = H_0 \psi. \rightarrow a(t;\epsilon) = a(0;\epsilon) e^{\epsilon t / i\hbar}.$$

$$\rightarrow \psi = \int d\epsilon a(0;\epsilon) e^{\epsilon t / i\hbar} |\epsilon\rangle. \rightarrow \partial_t |a(0;\epsilon) e^{\epsilon t / i\hbar}|^2 = 0. \rightarrow \text{"nothing quantum state transition"}$$

(2) "Unique Energy Eigen State Realization of H_0 " :

$$\{\text{stationarity} \Leftrightarrow \Delta t = \infty\}. \rightarrow \Delta E = 0. \text{ "unique energy eigen state realization of } H_0 \text{"}$$

In the below, it is also proved that degeneration in energy state is discriminated by other quantum numbers

②: Hermite Hamiltonian's Complete Causality and Markovian of General Quantum Process :

(1) "Realization of unique eigen state of maximum observable P_r of H_0 " : ($\mathbb{F} : H_0 |j\rangle = \epsilon^j |j\rangle$).

Maximum observable of $H_0 \equiv \{ P_r \mid [P_r, H_0] = 0 ; r = 1, 2, \dots, M \}$. $\dots \dots (1)-1$.

Now we prove $\Delta P_r = 0$ due to $\infty = \Delta Q_r, \Delta P_r = \hbar / \Delta Q_r$, where Q_r is canonical conjugate of P_r .

$$[Q_r, P_r] = i\hbar 1. \rightarrow i\hbar \partial_t Q_r = [Q_r, H_0] \neq 0. \dots \dots (1)-2$$

Such Q_r is time dependent in $\Delta t = \infty$ of H_0 . Hence if $q_r = q_r(t)$ is definite, $t = q_r^{-1}(q_r)$ must be also definite. This contradicts with indefiniteness of time as $\Delta t = \infty$ in H_0 system. Hence $\Delta Q_r = \infty$.

$$\{\text{stationarity} \Leftrightarrow \Delta t = \infty\}. \rightarrow \Delta P_r = 0. \text{ "unique eigen state realization of } H_0 \text{'s maximum observable } P_r \text{"}$$

$$i\hbar \partial_t \Psi = H_0 \Psi. \Leftrightarrow \{ \Psi = | \epsilon^j ; p_1^j, p_2^j, \dots, p_M^j \rangle \text{ is unique eigen state under hermitian Hamiltonian } H_0 \}$$

(2) Mathematical Singularity of Time Dependent Hamiltonian $\equiv H_s(t)$:

As the logical negation, Hamiltonian of causing state transition must be non-hermite. Then non-hermiteness must be non-observability of energy, that is $\Delta E = \infty$, which causes $\Delta t = 0$ for non-hermite Hamiltonian. In anyway such Hamiltonian never can be mathematically regular (non-analytical = singular). Hence the phenomena of $H_s(t)$ (\equiv non-hermite singular Hamiltonian) never can be causal. As its realities, observed quantum transitions are instantaneous phenomena ($\Delta t = 0$) with certain probability (\equiv non-causality).

Thus $H_s(t)$ shall cause quantum transition with certain probability in its duration $\Delta t = 0$ on time axis.

Then when time spot phenomenon of $H_s(t)$ shall realize becomes serious problem. $\rightarrow (4)$.

example 1) $H_{GF} = \int dx^3 [g \bar{\psi}(x) \gamma^\mu A_\mu(x) G_\alpha \psi(x)]$ is called minimal gauge interaction of spinor field ψ and gauge field $A_\mu(x)$ which are mathematically called "hyper function by Satch (or distribution by Schwarz)" due to Dirac's delta function in field commutation relation. Then $\psi \times A_\mu$ product with common field parameter $\equiv x$ never mathematically defined in general. That is breakdown of "regularity". As is well-known, physics of H_{GF} is probabilistical on its reactions in perturbation integral.

(3) Markovian Nature of General Quantum Process :

As the logic, general quantum process must be time serie of alternating realization of $\{ H_0 \ \& \ H_s(t) \}$. Hence realizable sample process must be also series of instantaneous probabilistical transitions among H_0 's eigen states. This is nothing without Markov Process of quantum one. \rightarrow Quantum Stochastic Mechanics.

③: Stochastic Hamiltonian and its Statistics (Evolution Theorem (or Principle) by Energy Fluctuation) :

(1) Fundamental idea = correlation function of $\Psi(t)$ as measure of initial state decay caused by $H_s(t)$.

As was mentioned, time dependent Hamiltonian $H_R(t)$ becomes random alternating realization of $H_0 \ \& \ H_s(t)$. Now we shall establish statistics of $H_s(t)$, of which effects should reflect on statistics on decay of initial state $\Psi(t) = \int dq a_q(t) |q\rangle$ to $\Psi(t + \Delta t)$, where $|a_q(t)|^2 \equiv \omega_q(t)$ is state density of $|q\rangle$. Because the variation of $a_q(t)$ is caused from realization of $H_s(t)$. The measure becomes correlation function between $\Psi(t)$ and $\Psi(t + \Delta t)$. Now we introduce "Modified Winer Kintchin Theorem" in correlation function and spectrum density. The modification is taken for the time integration interval which is finite in macroscopic view point, however it is sufficient to be infinite in microscopic quantum view point.

② Modified Winer Kintchin Theorem and Evolution Theorem by Energy Fluctuation (\equiv ETEF) :

(1) Semi-macroscopic integral time duration : $T(t) \equiv [t - \frac{1}{2}\Delta T, t + \frac{1}{2}\Delta T]$, where $\Delta T \gg \Delta t = \hbar / \Delta \epsilon$.

(2) Fourier component of $\Psi(t)$: $c(\epsilon; t) \equiv (2\pi\hbar)^{-1/2} \int_{T(t)} du |\Psi(u)\rangle \exp(\epsilon u / i\hbar)$.

(3) State density : $\omega(\epsilon; t) \equiv \langle c(\epsilon; t) | c(\epsilon; t) \rangle / \Delta T = \Delta T^{-1} (2\pi\hbar)^{-1} \int_{T(t)} du \int_{T(t)} dv \langle \Psi(u) | \Psi(v) \rangle \exp[-\epsilon(u-v) / i\hbar]$.

(4) Inverse Fourier : $\int_{-\infty}^{\infty} d\epsilon \exp(-\epsilon \Delta u / i\hbar) \omega(\epsilon; t) =$

$$\text{transform} = \Delta T^{-1} (2\pi\hbar)^{-1} \int_{T(t)} du \int_{T(t)} dv \langle \Psi(u) | \Psi(v) \rangle \int_{-\infty}^{\infty} d\epsilon \exp(-\epsilon(\Delta u + u - v) / i\hbar)$$

$$= \Delta T^{-1} \int_{T(t)} du \int_{T(t)} dv \langle \Psi(u) | \Psi(v) \rangle \delta(\Delta u + u - v) = \Delta T^{-1} \int_{T(t)} du \langle \Psi(u) | \Psi(u + \Delta u) \rangle \equiv \Upsilon(\Delta u; t).$$

(5) $\Upsilon(\Delta u; t) = \int_{-\infty}^{\infty} d\epsilon \exp(-\epsilon \Delta u / i\hbar) \omega(\epsilon; t) = 1 + i \langle \epsilon \rangle \Delta u / \hbar - \frac{1}{2} \langle \epsilon^2 \rangle \Delta u^2 / \hbar^2 + \dots$

(6) $|\Upsilon(\Delta u; t)| = 1 - \Delta \epsilon \Delta u / \hbar + \dots =$ non-decaying probability of initial $\Psi(t)$ by $H_s(t)$.

where $\Delta \epsilon(t) = \sqrt{\langle \epsilon^2 \rangle - \langle \epsilon \rangle^2}$ is statistical deviation of energy by $\omega(\epsilon; t)$. $\langle \epsilon^n \rangle = \int d\epsilon \cdot \epsilon^n \omega(\epsilon; t)$.

— EVOLUTION THEOREM* BY ENERGY FLUCTUATION — (*): or principle)

- (7): $\Theta(t) = \Delta \epsilon(t)/\hbar$: probability rate of generating $H_s(t)$ causing Ψ 's decay (state transition).
 (8): $\Delta \epsilon^2 = \langle \epsilon^2 \rangle - \langle \epsilon \rangle^2$: energy deviation in state density $\omega(\epsilon; t)$.
 (9): $\Delta \epsilon \Delta t = \hbar$: uncertainty theorem for time and energy.
 (10): $\Delta t = \hbar/\Delta \epsilon$: average duration of H_0 , or average rate time of generating $H_s(t)$.

4: Establishing Quantum Stochastic Mechanics and the Problems of General Isolated Closed Systems:

- ① Markov process $\Leftrightarrow \partial_t \omega_j(t) = \sum_k \Gamma_{jk}(t) \omega_k(t) - \sum_k \Gamma_{kj}(t) \omega_j(t)$. < Master equation >
 ② Transition probability rate $\Gamma_{jk}(t) = \Theta(t) T_{jk} = [\Delta \epsilon(t)/\hbar] T_{jk}$.
 (1) 1st order transition probability by $H_s(t)$: $T_{jk} = | \langle i | \int_0^t dt \langle j | H_s(t) | k \rangle |^2$. (\mathcal{E} : U is not known)
 To tell frankly, concrete details of T_{jk} has not been surveyed yet. According to perturbation integral method in QFT, coordinate x dependent term in the integral is $\int_0^t dx^4 \exp[(p^t_\mu - p^i_\mu) x_\mu / \hbar] = (2\pi\hbar)^4 \delta(p^t_\mu - p^i_\mu)$, where $U = [-\infty; \infty]$. This represents 4-dimensional momentum conservation law for $|i\rangle \rightarrow |f\rangle$ in U.
 Therefore T_{jk} dominant term may be other constant term such as "function of momentum space variables".
 (2) See ⑥(6). There is a possibility of deriving T from final equilibrium state $\equiv \omega^\infty$, where $T \omega^\infty = \omega^\infty$.
 (3) In view point of Goedel theorems, there must be algorithm on T_{jk} . Because a probabilistical theorem itself is a unique deterministic description. There never be breaking down of uniqueness of theory.

③ $\partial_t \omega_j(t) = \langle \Delta \epsilon(t)/\hbar \rangle \sum_k [T_{jk} - \delta_{jk}] \omega_k(t)$. < reaction rate $\equiv 1/\Delta t(t) = \Delta \epsilon(t)/\hbar$: uncertainty of $\Delta \epsilon$ & Δt >
 $\partial_t \omega(t) = \langle \Delta \epsilon(t)/\hbar \rangle [T - 1] \omega(t)$.

The equation yields following important realities on isolated closed thermo-dynamical system.

- ④ **Irreversibility of the Eqn:** proof) $t \rightarrow (-t)$ is not invariant for form of the equation.
 ⑤ **Entropy Increasing Law:** proof) $S \equiv k_B \sum \omega_j \ln(1/\omega_j)$.
 $\partial_t S(t) = \frac{1}{2} k_B \sum_{j,k} (T_{kj} \omega_j - T_{jk} \omega_k) (\ln \omega_j - \ln \omega_k) \geq 0$, $\leftarrow (T_{kj} = T_{jk})$.

⑥ General Relaxation Process Solution in Isolated Closed System:

- (1) Markov chain expansion solution: $\omega = \sum_{n=0}^{\infty} R_n(t) T^n \omega_0$. < $\omega_0 \equiv \omega_0(0)$: initial state density >
 (2) $R_n(t)$ is the probability of realization of "n th order reaction at time=t".
 $R_n(t) \geq 0$, $1 = \sum_{n=0}^{\infty} R_n(t)$. \Rightarrow (5): domino-propagation of peak value $R_n(n=0-1-2-3-\dots)$.
 (3) $0 = \partial_t R_0 + \Theta R_0 \Rightarrow R_0 = \exp[-\int_0^t du \Theta(u)]$.
 $0 = \partial_t R_{n+1} + \Theta R_{n+1} = \Theta R_n \Rightarrow R_n = \int_0^t du \Theta(u) R_{n-1}(u) \exp[-\int_0^u ds \Theta(s)]$. ($n=1, 2, 3, \dots$).
 (4) $R_0(t)$ is monotonous decreasing function from 1 at $t=0$.
 (5) $\partial_t R_{n+1} = \Theta (R_n - R_{n+1})$: < note that $\Theta > 0$ >
 R_{n+1} shall increase from zero and to have single maximum point at $R_n(\searrow) = R_{n+1}(\nearrow)$, then becomes monotonously decreasing function toward zero due to $(R_n - R_{n+1}) < 0$.
 (6) $R_n(\infty) \rightarrow +0 \Rightarrow T \omega(\infty) = \omega(\infty) \Rightarrow$ equilibrium state = general relaxation process.
 \mathcal{E} : In (6), there is a possibility of deriving T from well known equilibrium state $= \omega(\infty) \equiv \omega^\infty$.
 Because $0 = \sum_k T^{(j)}_{jk} \omega^\infty_k$, hence vector ω^∞ is orthogonal to $T^{(j)}$, where $T^{(j)}_{jk} \equiv T_{jk}$.

5: The Problems of General Opened Systems with Thermo-Chemical External Flows:

- ① State transition is caused also by in & out thermo-chemical flow at boundary of the system, which is equivalent to being another singular Hamiltonian $\equiv J_s(t)$ like as $H_s(t)$. Then we assume probabilistical exclusiveness as follows. $\Gamma_{jk}(t) \equiv \Gamma^i_{jk}(t) + \Gamma^e_{jk}(t) = \Theta(t) T_{jk} + \Lambda(t) L_{jk}$. $\dots \dots$ (1)

(2) $\partial_t \omega_j(t) = \langle \Delta \epsilon_i(t)/\hbar \rangle \sum_k [T_{jk} - \delta_{jk}] \omega_k(t) + \langle \Delta \epsilon_L(t)/\hbar \rangle \sum_k [L_{jk} - \delta_{jk}] \omega_k(t)$. \dots opened system eqn.
 (3) $\Delta \epsilon_L(t) = \sqrt{\sum_{j,k} L_{jk}(t) \omega_k(t) (\epsilon_j - \langle \epsilon \rangle)^2}$. : Energy deviation caused by flowing transition L_{jk} .

② Simple model solution \equiv virtual stationary flow and heart beating solution: $j_0 \equiv \Lambda(t) [L - 1] \omega(t)$.

- (1) $\partial_t \omega = \Theta [T - 1] \omega + j_0$; $\omega \equiv \sum_{n=0}^{\infty} R_n(t) T^n \omega_0 + \sum_{m=0}^{\infty} F_m(t) T^m j_0$.
 $0 = \partial_t R_0 + \Theta R_0$; $0 = \partial_t R_{n+1} - \Theta (R_n - R_{n+1})$; $0 = (\partial_t F_0 + \Theta F_0 - 1)$; $0 = \partial_t F_m - \Theta (F_{m-1} - F_m)$.
 (2) R_n behaves the same as one in isolated closed system. $\Rightarrow R_n(\infty) = 0$.
 (3) $0 = \partial_t F_0 + \Theta F_0 - 1 \Rightarrow F_0(t) = \exp[-\int_0^t ds \Theta(s)] \int_0^t du \exp[\int_0^u ds \Theta(s)] \geq 0$.
 $0 = \partial_t F_m - \Theta (F_{m-1} - F_m) \Rightarrow F_m(t) = \int_0^t du \Theta(u) F_{m-1}(u) \exp[-\int_0^u ds \Theta(s)] \geq 0$.
 (4) $\partial_t F_m(\infty) = 0 \Rightarrow \{F_m(\infty) = 1/\Theta(\infty) = \Delta t(\infty)\}$; (2) $\Rightarrow 0 = \partial_t \omega(\infty) = \Theta(\infty) [T - 1] \omega(\infty) + j_0$.
 (5) $\omega(\infty) = T \omega(\infty) + \Delta t(\infty) j_0$: < solution of heart beating with stationary flow >.
 $\omega(\infty) = \Delta t(\infty) [1 - T]^{-1} j_0$: < equilibrium state determined by flow j_0 and T >.

-SUPPLEMENT-

① The Fundamental Axioms of Quantum Mechanics (as trial rough proposal) :

A0 : any observable physical variable is represented by hermitian operator A and observed value is their eigen value a_p , with "eigen state function" $\psi = |a_p\rangle$. $\rightarrow A |a_p\rangle = a_p |a_p\rangle$. < eigen equation >

A1 : classical (\equiv non-quantized) mechanical system of Lagrangean $L(Q_i, \dot{Q}_i, Q_j)$ is quantized by Canonical Quantization Principle (\equiv CQP) on physical variables as hermite operator as follows.

$$P_j \equiv \partial L / \partial (\dot{Q}_j) \rightarrow [Q_j, P_k] \equiv Q_j P_k - P_k Q_j = i\hbar \delta_{jk} 1.$$

$$[Q_j, Q_k] = [P_j, P_k] = 0.$$

E7 : The principle determining classical field Lagrangean are (1) global Lorentz covariance, (2) localized Lorentz one. (1) is for free spinor field Lagrangean, and (2) is for quantum gravitational Lagrangean of unified field. They are unifiedly called "Transform Invariance Principle (\equiv TIP)".

Theorem1 : Observed value is expressed as $a_p = \langle a_p | A | a_p \rangle$.

T2 : Commutable observables has common eigen function and "enables simultaneous observation".

proof) $0 \equiv [A, B] \rightarrow B = F(A) \rightarrow B |a_p\rangle = F(a_p) |a_p\rangle$.

E7; If $A = H_0$, B is MO \equiv maximum observable.

T3 : Eigen functions becomes complete ortho-normal function set.

Any function $|c_q\rangle$ can be expressed its expansion form as $|c_q\rangle = \int dq' \cdot u_{pq'} |a_p\rangle$.

$\rightarrow 1 \equiv \langle a_p | a_p \rangle = \int dq' dq'' \cdot u_{pq'}^* \cdot u_{pq''} \langle c_q | c_q \rangle = \int dq' dq'' \cdot u_{pq'}^* \cdot u_{pq''} \delta(q' - q'') = \int dq' |u_{pq'}|^2$.

Hence $\{|u_{pq'}|^2\}$ has "probability feature".

T4 : Simultaneous observation on state $|a_p\rangle$ by non-commutable observable C never can yield unique p "deterministic results", but yield statistical results < Breakdown of causality in non-MO observation >

r $C |c_q\rangle = c_q |c_q\rangle \rightarrow$ expansion theorem $\rightarrow |a_p\rangle = \int dq' \cdot v_{pq'} |c_q\rangle$.

o \rightarrow Observed value (T1) $= \langle a_p | C | a_p \rangle = \int dq' \int dq'' \cdot v_{pq'} \cdot v_{pq''}^* \cdot c_q \langle c_q | c_q \rangle = \int dq' |v_{pq'}|^2 c_q$.

f From axiom A0, each sample observation value must be eigen value c_q . However observed value (T1) has averaged value form with probability density $|v_{pq'}|^2$. Therefore causality uniqueness of non-commutable observation is broken down and is to ensure "statistical interpretation".

E7 : See Goedel Incompleteness theorem.

In this way, in our quantum mechanics, introducing probability is not an axiom, but is a theorem.

T5 : Non-commutable observation is a irreversible reaction on initial state $|a_p\rangle$ into $|c_q\rangle$.

The irreversibility is evident by entropy increasing $S = k_b \int dq' |v_{pq'}| \ln(1/|v_{pq'}|) > 0$.

E7 : $\psi = \psi(p; x)$ is assumed to be a momentum eigen function. If we try particles position observation, then

$\psi = \int dx' \psi(p; x') \delta(x - x') \rightarrow \delta(x - x')$ of particle position x' with the probability density $= |\psi(p; x')|^2$

In the reality, $\psi \rightarrow \delta$ is a reaction by external injection of test particle for observation.

EXT6 : Relich-Dixmier theorem : such variables $[P_j, Q_j] = i\hbar 1$ are transformed into $\{P_j = -i\hbar \partial / \partial x_j, Q_j = x_j\}$ by certain unitary transform in general.

T7^o) : Deriving Schroedinger equation from CQP :

As for time "t", its has two canonical conjugate variable as $\{H_0 \equiv$ Hamiltonian, $i\hbar \partial_t \equiv$ time derivative $\}$, due to dimension analysis in CQP.

Because $x_\mu = (ict, x_1, x_2, x_3)$ and $p_\mu = (iE/c, p_1, p_2, p_3) = -i\hbar \partial_\mu$ (EXT6 and external theory of relativity).

Hence $E = (c/i) (-i\hbar \partial / \partial x_0) = (c/i) (-i\hbar \partial / ic \partial t) = i\hbar \partial / \partial t$. Then $H_0 = i\hbar \partial_t$ for any functions? Absolutely no!.

Unique possibility is $H_0 = i\hbar \partial_t$ for certain function Ψ . That is $i\hbar \partial_t \Psi = H_0 \Psi$.

E7 : Certainly "theory of relativity is out of this axiom system", however in unified principles with TIP, it may can be closed?.

T8 : Uncertainty Theorem for Canonical Conjugate Variable $[P_j, Q_j] = i\hbar 1$:

Then we can prove following inequality in which $\Delta P_j, \Delta Q_j$ are statistical deviation.

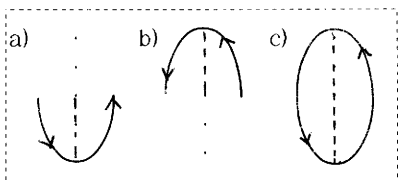
$$\Delta P_j \cdot \Delta Q_j \geq \frac{1}{2} \hbar.$$

T9 : As for time and energy, it is rather complicated to derive. The result is $\Delta E \cdot \Delta t = \hbar$.

② -Diffraction Pattern Forming by Single Electron Beam's Accumulation through Two Slits-

After all, single electron of constant momentum can go through two slits simultaneously by "random instantaneous space transportation of electron" through higher order vacuum polarization reactions. A free running elementary particle never be free, but is also a consequence with vacuum field reactions. Because quantum vacuum field is supremely filled with vacuum polarization reactions at anywhere.

(1) A quantum vacuum field is not empty but is filled with vacuum polarization reactions of any kind of spinor elementary particles pairs in any space and any time. Such reactions are caused from H_{GF} .



$$H_{GF} = g\chi\bar{\psi}\gamma^\mu A_\mu G_\alpha\psi \text{ (minimal gauge interaction)}$$

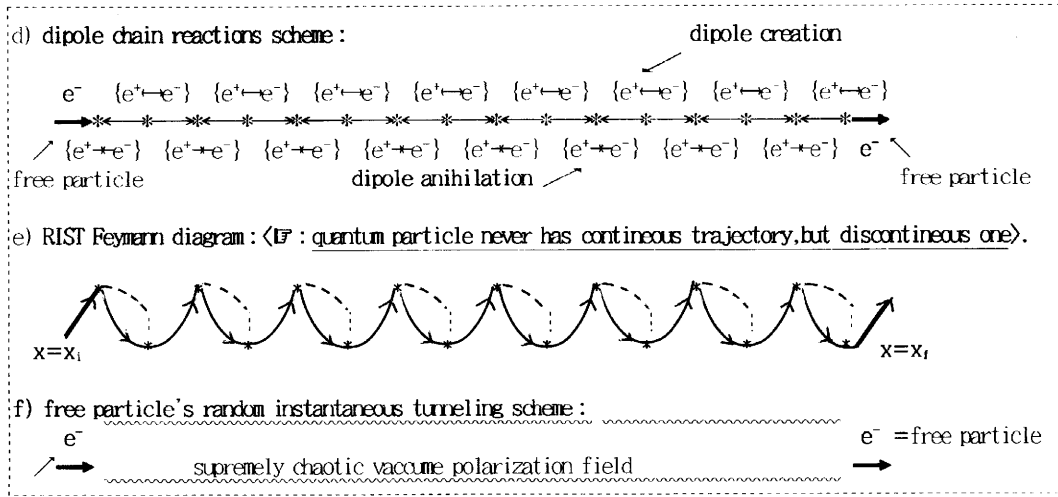
$$\begin{aligned} \bar{\psi}(x) &= \sum_s \int d^3p \{ b(\mathbf{p},s) \bar{v}(\mathbf{p},s) e^{-ipx/\hbar} + a^\dagger(\mathbf{p},s) \bar{u}(\mathbf{p},s) e^{ipx/\hbar} \}, \\ \psi(x) &= \sum_s \int d^3p \{ a(\mathbf{p},s) u(\mathbf{p},s) e^{-ipx/\hbar} + b^\dagger(\mathbf{p},s) v(\mathbf{p},s) e^{ipx/\hbar} \}, \\ A_\mu^\alpha(x) &= \sum_\lambda \int d^4q \{ c^\alpha(\mathbf{q},\lambda) \epsilon_\mu(\mathbf{q},\lambda) e^{-iqx/\hbar} + c^{\alpha\dagger}(\mathbf{q},\lambda) \epsilon_\mu(\mathbf{q},\lambda) e^{iqx/\hbar} \}. \end{aligned}$$

Consequently 8 kind of 1st order reactions are derived from H_{GF} . Any kind of higher order quantum field reactions are time series of each 8 kind of fundamental reactions. Especially vacuum polarization creation a) and vacuum polarization annihilation b) are fundamental. c) is sequential reaction of a)→b) as closed vacuum polarization as 2nd order reaction. Different reactions exist from that of perturbation theory.

(2) An elementary particle never can run through such vacuum field without any collision with particles of vacuum polarizations at anywhere! Usually a collision of particle becomes a reaction.

In former interpretation, free running elementary particle is considered without such fundamental reactions. Telling fact, being free running particle in non-localized space is caused from higher order of such ones. Because an elementary particle never can run through such vacuum field without collision with particles of vacuum polarizations being supremely filled at anywhere! The reaction is as follows.

(3) Elementary Particle's Instantaneous Random Space Transportation (≡IRST=discontinuous free running) Through Vacuum Polarization Dipole Chain Reactions=Free Particle's Tunneling Mechanism:



(4) Thus you will see that an elementary particle which is certainly not being with zero volume simultaneously never be local being, but is non-localized being with random instantaneous space transportation (≡RIST) through chaotic vacuum polarization field with "definite memory of constant momentum" as a plane quantum wave function $\psi = \exp(ipx/\hbar)$. ψ is of course an eigen function of momentum observable which is commutable with H_0 . In addition to tell, such RIST is also possible for "any kind of complex particles" due to being of nucleon dipole forming reaction with FP ghost in general gauge field theory (author).

(5) Quantum wave eigen function $\psi(x)$ of position variable is a reality. The statistical interpretation of position observation on particle is due to such RIST of particle. Hence, in ψ (=state without observation on it), a particle can be two slits simultaneously by RIST which itself is of course non-observable. Therefore RIST never contradict with law of upper limit of velocity of light in observable physics. In addition to tell, single electron can go through "any N pieces of slits" simultaneously.